

Last time: Problems in 2 or more dimensions

$$\text{2-D or 3-D box } \hat{H} = \frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right)$$

$$= \hat{H}_x + \hat{H}_y + \hat{H}_z \quad \{ \text{sum}$$

$$\Rightarrow \psi(x, y, z) = \psi(x) \cdot \psi(y) \cdot \psi(z) \quad \{ \text{product}$$

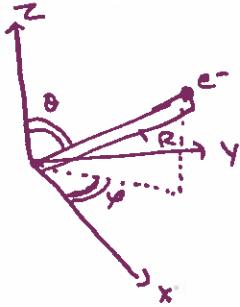
$$\hat{H} = \hat{H}_{\text{rot}}(\theta) + \hat{H}_{\text{vib}}(r) \quad \{ \text{sum}$$

$$\psi(\theta, r) = \psi(\theta) \cdot \psi(r) \quad \{ \text{product}$$

ex- rotating -
vibrating
molecule

Convenient for more
advanced types of particles

To day: rotation in 3D



θ and ϕ may be
flipper in mathematics
bodies

$$\hat{H} = \frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) + V(x, y, z)$$

↓ change to polar coordinates

$$r = \sqrt{x^2 + y^2 + z^2}$$

$$\psi = \tan^{-1}(y/x)$$

$$\theta = \tan^{-1}(\sqrt{x^2 + y^2}/z)$$

$$\hat{H}(\psi, \theta, r) = \frac{-\hbar^2}{2mr^2} \left(\frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \psi^2} + \frac{\partial^2}{\partial \theta^2} + \frac{\cos \theta}{\sin \theta} \frac{\partial}{\partial \theta} \right) +$$

$$\left\{ -\frac{\hbar^2}{2m} \left(\frac{\partial^2}{\partial r^2} + \frac{1}{r} \frac{\partial}{\partial r} \right) + V(r) \right\}$$

looks the
same as the
rotating-vibrating
molecule

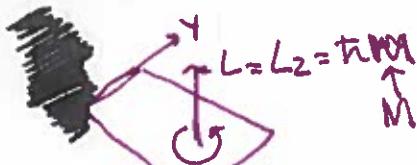
interested in
Coloumb potential
(only depends on r)

$$\hat{H} = \hat{H}_{\text{rot}}(\theta, \phi) + \hat{H}(r)$$

$$\psi = \psi(\theta, \phi) \cdot \psi_r(r)$$

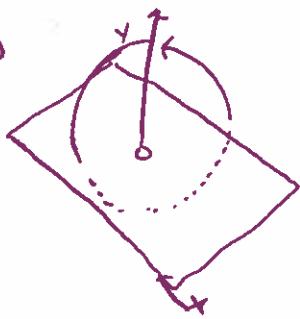
$$\text{As in 2D, } \hat{H}_{\text{rot}} = \frac{\hbar^2}{2mR^2}$$

Difference between 2D and 3D rotation



$$\hat{H}_{\text{rot}} = \frac{\hbar^2}{2mR^2} = \frac{\hbar^2}{2mR^2} \frac{\partial^2}{\partial \psi^2} \quad \psi_{\text{rot}}(\psi) \sim e^{iM\psi}, L_z = \hbar M$$

3D



Angular Momentum
Can point anywhere
in 3D

In the case above $L_z = \hbar M = 0$

But " $l \neq 0$ "
(ℓ = total angular momentum)

The solutions for $\hat{H}_{\text{rot}} + \hat{Y}_{\text{rot}}(\theta, \varphi) = E_{\text{tot}} Y_{lm}(\theta, \varphi)$ can be derived using raising and lowering operators, but ... not going through that again

$$\frac{\hbar^2}{2mr^2} Y_{lm}(\theta, \varphi) = \frac{\hbar^2 l(l+1)}{2mr^2} \stackrel{E_{\text{tot}}}{\leftarrow} Y_{lm}(\theta, \varphi) \quad l=0, 1, \dots$$

We have two coordinates $\theta, \varphi \Rightarrow$ two quantum #s

l and M
 \uparrow
total angular momentum
 $l=0, 1, 2$

z-axis angular momentum $M=0, \pm 1, \pm 2$

$$|M| \leq l$$

The length of a vector can't be longer than its z-axis component

Lecture on Monday \rightarrow what these wavefunctions look like