

## Exam 1 Solution

Useful numbers and plots you may need are at the end.

1. (18 pts) **Calculate** the wavelength of light that will be absorbed by the molecule propene (it has a C-C=C backbone) by approximating its electronic structure by an “electron in a box.”

This one was similar to problem 2.4 in the book.

a. (2) With  $d(\text{C-C}) = 1.54 \text{ \AA}$  and  $d(\text{C=C}) = 1.35 \text{ \AA}$ , **what is** the length  $L$  of the box, assuming the electron can move through the whole carbon skeleton?

Solution:  $L = 2.89 \text{ \AA}$  (2 points)

b. (3+3) Of the three  $\pi$  electrons in propene, the first two will fill the  $n=1$  energy level of the box, and the third electron, which will be excited by the light, will be in the  $n=2$  level or “HOMO” (highest occupied molecular orbital). **Write down** the particle-in-a-box energy formula for the third electron in terms of  $n$ ,  $h$ ,  $m_e$  and  $L$  and **calculate** the energy of the  $n=2$  electron in Joules.

$$\text{Solution: } E_n = \frac{n^2 h^2}{8mL^2} \quad (3 \text{ points})$$

$$\begin{aligned} &= \frac{(2)^2 (6.626 \cdot 10^{-34} \text{ Js})^2}{8(9.1 \cdot 10^{-31} \text{ kg})(2.89 \cdot 10^{-10} \text{ m})^2} \\ &= 2.88 \cdot 10^{-18} \text{ J} \quad (3 \text{ points}) \end{aligned}$$

c. (3+2) When light is absorbed, the electron goes to the  $n = 3$  level. **Write down** the formula for the energy difference between the  $n = 3$  and  $n = 2$  states, and **calculate** the energy difference in Joules.

$$\text{Solution: } \Delta E = (3^2 - 2^2) \frac{h^2}{8mL^2} \quad (3 \text{ points})$$

$$\begin{aligned} \Delta E &= (3^2 - 2^2) \frac{(6.626 \cdot 10^{-34} \text{ Js})^2}{8(9.1 \cdot 10^{-31} \text{ kg})(2.89 \cdot 10^{-10} \text{ m})^2} \quad (2 \text{ points}) \\ &= 3.61 \cdot 10^{-18} \text{ J} \end{aligned}$$

d. (3+2) Finally, use Planck’s law to convert the energy to a **wavelength in nm**. What part of the **electromagnetic spectrum** at the end of the exam is it?

$$\begin{aligned} \text{Solution: } 3.61 \cdot 10^{-18} \text{ J} &= \frac{hc}{\lambda} = \frac{(6.626 \cdot 10^{-34} \text{ Js})(3.0 \cdot 10^8 \text{ m/s})}{\lambda} \quad (3 \text{ points}) \\ \lambda &= 6.2 \cdot 10^{-8} \text{ m} = 62 \text{ nm} \end{aligned}$$

This is in the extended UV range (2 points)

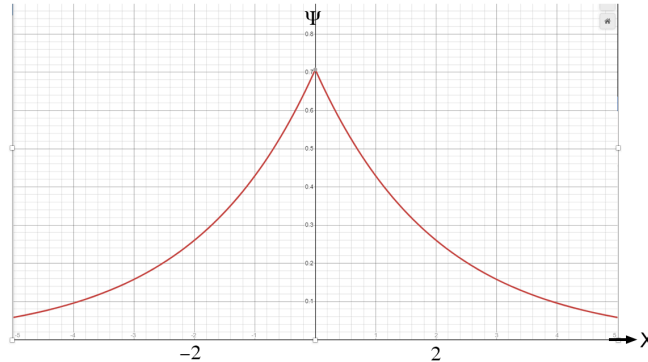
Note that experiments show that propene absorbs in the mid-UV, at longer wavelengths. The reason is that the Coulomb potential is a ‘softer’ box than the particle in a box. However, the box approximation gets better the longer the conjugated chain is, as the ‘edge effects’ are smaller in those cases.

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2. (15 pts) The normalized wavefunction  $\Psi(x) = \left(\frac{1}{\sqrt{2}}\right) \exp[-|x|/2]$  is also known as a “1s orbital” for an electron whose nucleus sits at  $x=0$ . Here  $x$  is in Ångstrom units.

a. (4+2+2) **Sketch** the wavefunction roughly between  $x=-5$  to  $5$  Å, **labeling** the axes “ $x$ ” and “ $\Psi$ ” and **indicating** roughly where  $x=\pm 2$  lies on the  $x$ -axis.

Solution: This one was similar to problem 1.8 in the book and H12, #2



b. (2+3+2) Calculate the probability of finding the electron at a distance between 3 and 3.1 Å from the nucleus by first **writing down the integral** with correct integration limits, then **evaluating** the integral, and **plug in** the numbers to get the probability.

Solution:

$$\begin{aligned} P &= \int_3^{3.1} |\varphi(x)|^2 dx \\ &= \int_3^{3.1} \left| \left(\frac{1}{\sqrt{2}}\right) \exp[-|x|/2] \right|^2 dx \\ &= \frac{1}{2} \int_3^{3.1} e^{-|x|} dx = \frac{-1}{2} e^{-x} \Big|_3^{3.1} \\ &\text{Integrate from 3 to 3.1} \approx 0.00237 \end{aligned}$$

The actual probability is twice that because electrons at  $-3$  to  $-3.1$  Å are also at a distance of between 3 and 3.1 Å from the nucleus. However, we'll give full credit for either answer if you overlooked that!

3. (10 pts) Someone plays the kettle drum.

a. (3+3) If they play on the kettle drum a “low “A” note at 110 Hz for 50 milliseconds (0.05 s), what is the uncertainty in the pitch  $\Delta v$  as a percentage of the frequency  $v$ ? **Formula** and **value**.

Solution: This one was similar to H5 #1

$$\Delta v \Delta t = \frac{1}{4\pi} \quad [\text{Note } v = \omega/2\pi]$$

Therefore

$$\Delta v = \frac{1}{0.2\pi} \approx 1.59 \text{ Hz} \quad (3 \text{ points})$$

$$\frac{\Delta v}{v} \times 100\% \approx 1.45\% \text{ (3 points for the formulas, 3 for the values)}$$

b. (2+2) **Can they distinguish** the “A” from an “A#” at 116.5 Hz? In one sentence, why do you think orchestras use kettle drums, but not snare drums, to play bass melodies in symphonies?

Solution: Yes. (2 points)

Kettle drums sound long enough that one can make out the pitch reasonably well to play a bass melody; snare drums don't. (2 points)

4. (18 pts) If a quantum particle is in state  $\Psi(x, t)$ , let us show that the average value of the energy from many measurements is given by

$$\langle E \rangle = \bar{E} = \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{H} \Psi(x, t),$$

a. (3+3) **Write down** the formula for  $P(E = E_n) = ?$  from postulate 4, calling the energy eigenvalues  $E_n$  and the eigenfunctions  $\varphi_n(x)$ . Then recall from the \$ bills homework that  $\bar{A} = \sum a_n P(A = a_n)$  for any observable A. **Combine these two equations** to write an expression for  $\langle E \rangle$  in terms of the eigenvalues and eigenfunctions of  $\hat{H}$ .

b. (3) Now **write down** the same expression for  $\langle E \rangle$  again, but multiply out the  $|\cdot|^2$  square modulus explicitly. Something like

$$\int_{-\infty}^{\infty} dx \Psi^*(x, t) \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t)$$

should appear in your expression (the integral times its complex conjugate). It's always a good idea to give your integration variables different names (here  $x$  and  $x'$ ) when you have several integrals in an expression.

c. (3) Your expression for  $\langle E \rangle$  has an “ $E_n$ ” outside the first integral. Since  $E_n$  is a constant, you can stick it inside the integral in front of  $\varphi_n(x)$ , and play the reverse of a trick we have done in class several times: since  $\hat{H}\varphi_n(x) = E_n\varphi_n(x)$ , you can replace “ $E_n$ ” **by what** in the integral? **Write down** the formula for  $\langle E \rangle$  again, this time without  $E_n$  in it.

d. (3+3) Remember that any function can be expanded as a linear combination of a complete set of eigenfunctions. For example,  $\Psi(x, t) = \sum c_m(t) \varphi_n(x)$ , where  $c_m(t) = \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t)$  is the “overlap integral” between  $\varphi_n(x)$  and  $\Psi(x, t)$ . Look at your expression from part c., and replace one of the integrals by  $c_m$ , and **write down** the resulting formula for  $\langle E \rangle = \bar{E}$ . Finally **do the sum**  $\sum c_m \varphi_n(x)$ , which equals  $\Psi(x, t)$ , and write down your final expression for  $\langle E \rangle = \bar{E}$ .

Congrats, in 1926 you would have won a Nobel Prize!

Solution: Similar to Practice Exam Fall 2016, Q. 5, if you substitute “H” instead of “A” and “ $E_n$ ” instead of “ $a_n$ ”

a. (3 pts)

$$P(E = E_n) = \sum_n E_n \left| \int_{-\infty}^{\infty} dx \varphi_n^*(x) \Psi(x, t) \right|^2$$

(3 pts)

$$\bar{E} = \sum_n E_n \left| \int_{-\infty}^{\infty} dx \varphi_n^*(x) \Psi(x, t) \right|^2$$

b. (3 pts)

$$\bar{E} = \sum_n E_n \int_{-\infty}^{\infty} dx \Psi^*(x, t) \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t)$$

c. (3 pts)

$$\begin{aligned} \bar{E} &= \sum_n \int_{-\infty}^{\infty} dx \Psi^*(x, t) E_n \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t) \\ &= \sum_n \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{H} \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t) \end{aligned}$$

d. (3+3 pts)

$$\begin{aligned} \bar{E} &= \sum_n \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{H} \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t) \\ &= \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{H} \sum_n \varphi_n(x) \cdot \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t) \\ &= \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{H} \sum_n \varphi_n(x) \cdot c_n(t) = \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{H} \Psi(x, t) \end{aligned}$$

This integral is the overlap coefficient  $c_n$

This last step we discussed in lecture 11: since the time-dependent and time-independent Schrödinger equations are equivalent (via Fourier transform), we can expand any  $\Psi(x, t)$  in terms of eigenfunctions of the energy  $\varphi_n(x)$ ; here we included the  $e^{-iE_n t/\hbar}$  in the phase factor  $c_n(t) = c_n e^{-iE_n t/\hbar}$  to keep the notation shorter, but of course, if you wrote that out in full, that's full credit.

Useful numbers:

1 atomic mass unit =  $1.66 \times 10^{-27}$  kg; mass of electron  $m_e = 9.109 \times 10^{-31}$  kg

Planck's constant  $h = 6.626 \times 10^{-34}$  J·s; note that  $\hbar = 2\pi h$  is about 6.28 times larger.

1 m = 100 cm; 1 Å = 0.1 nm = 100 pm

