

Lecture 30

k_{ij} metric van der Waals interaction exchange term

$$E_0 \approx \langle \hat{H} \rangle = \sum_i \langle i | \hat{h}_i | i \rangle + \frac{1}{2} \sum_{i,j \neq i} J_{ij} - \frac{1}{2} \sum_{i,j \neq i} k_{ij}$$

Coulomb repulsion

$$= \sum_i \langle i | \hat{h}_i | i \rangle + \frac{1}{2} \sum_{i,j \neq i} \langle ij | \frac{1}{r_{ij}} | ij \rangle - \frac{1}{2} \sum_{i,j \neq i} \langle ij | \frac{1}{r_{ij}} | ji \rangle$$

$$\phi_i(\vec{r}_i) = \sum_n c_n \psi_n(\vec{r}_n) \cdot \alpha_n \text{ or } \beta_n$$

include σ -orbitals in the calc.

6-3-6
 # of orbitals for core e-
 2 basis fn for valence e-'s
 36+16

gaussian

Gaussian Basis



$$\psi_0(\vec{r}_1, \vec{r}_2, \dots, \vec{r}_N) = \frac{1}{N!} \phi_1(\vec{r}_1) \dots \phi_N(\vec{r}_N)$$

How to systematically get the best C_m that gives the lowest EHF?

A: SCF turns on N -e- S.eq. into N 1e- S.eq.

Write EHF as

$$E_{HF} = \sum_i \langle i | \hat{h}_i | i \rangle + \frac{1}{2} \sum_{i,j \neq i} \langle ij | \frac{1}{r_{ij}} | ij \rangle - \frac{1}{2} \sum_{i,j \neq i} \langle ij | \frac{1}{r_{ij}} | ji \rangle$$

$$= \sum_i \langle i | \left[\hat{h}_i + \frac{1}{2} \sum_{j \neq i} J_{ij} - K_{ij} \right] | i \rangle = \sum_{i=1}^N \epsilon_i$$

$$\hat{J}_{ij} | i \rangle = \langle j | \frac{1}{r_{ij}} | j \rangle | i \rangle$$

$$K_{ij} | i \rangle = \langle j | \frac{1}{r_{ij}} | i \rangle | j \rangle$$

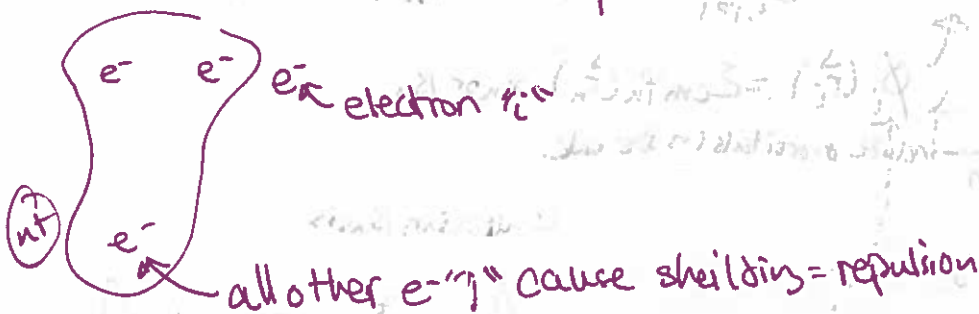
$$J_j(\vec{r}_i) \psi_i(\vec{r}_i) = \psi_i(\vec{r}_i) \int \frac{\psi_j^*(\vec{r}_j) \psi_j(\vec{r}_j)}{r_{ij}}$$

switches = exchanges the $|i\rangle$ and $|j\rangle$ functions

Now remove the $\sum_{i \neq j} <i|$ and solve the resulting N equations separately, each with its own $1e^-$ orbital energy ϵ_i

$$\Rightarrow (\hat{h}_i + \frac{1}{2} \sum_{j \neq i} \hat{J}_j(\vec{r}_i) - K_1(\vec{r}_i)) |i\rangle = \epsilon_i |i\rangle \text{ or } \hat{H}_i |i\rangle = \epsilon_i |i\rangle$$

has averaged over all other e-'s " $j \neq i$ "
but depends only on \vec{r}_i .



Proof: see N30b

1) Guess $|i\rangle = \phi_i = \sum c_m \phi_m(\vec{r}_i) | \pm \frac{1}{2} \rangle$

2) Calc. \hat{J}_j & K_j by doing integrals

such as $J_2(\vec{r}_1) = \int d\vec{r}_2 \frac{\phi_2^*(\vec{r}_2) \phi_2(\vec{r}_2) e^2}{4\pi\epsilon_0 r_{12}}$

$K_2(\vec{r}_1) = \int d\vec{r}_2 \frac{\phi_2^*(\vec{r}_2) \phi_1(\vec{r}_2) e^2}{4\pi\epsilon_0 r_{12}}$

3) Solve $\hat{H}_i |i\rangle = \epsilon_i |i\rangle$ to get new $|i\rangle$'s $\equiv \phi_i$

Upon convergence, you get the best possible determinant wavefn by letting

$$\psi_0 = \frac{1}{\sqrt{N!}} \begin{vmatrix} \phi_1(i_1) & \dots & \phi_n(i_1) \\ \phi_1(i_2) & \dots & \phi_n(i_2) \\ \vdots & & \vdots \\ \phi_1(i_N) & \dots & \phi_n(i_N) \end{vmatrix}$$

If not good enough, try diff e^- configurations $\psi_{n \neq 0}$

ground state ψ_0

ψ_2 doubly excited state ψ_2

$$\Rightarrow \frac{H}{2} = \begin{pmatrix} H_{00} & H_{02} & \dots \\ H_{20} & H_{22} & \dots \\ \vdots & \vdots & \ddots \end{pmatrix}$$

$H_{01} = 0!$