

Lecture 27

See review sheet

Localized orbitals: valence bonds or hybrid basis functions

Last time: ground state of $H_2 \approx \psi_{MO} = \frac{1}{\sqrt{2}} (|\sigma\rangle_1 + |\sigma\rangle_2) \{ \alpha_1\beta_2 - \alpha_2\beta_1 \}$

\uparrow \uparrow
 bonding MOs



$$\psi_{MO} = \frac{1}{\sqrt{2}} (|s_A\rangle_1 + |s_B\rangle_1) (|s_A\rangle_2 + |s_B\rangle_2) \cdot \text{spin}$$

$$= \frac{1}{\sqrt{2}} \left\{ \underbrace{|s_A\rangle_1 |s_A\rangle_2}_{H^+ + H^+} + \underbrace{|s_A\rangle_1 |s_B\rangle_2}_{H\cdot + H\cdot} + \underbrace{|s_B\rangle_1 |s_A\rangle_2}_{H\cdot + H\cdot} + \underbrace{|s_B\rangle_1 |s_B\rangle_2}_{H^+ + H^-} \right\}$$

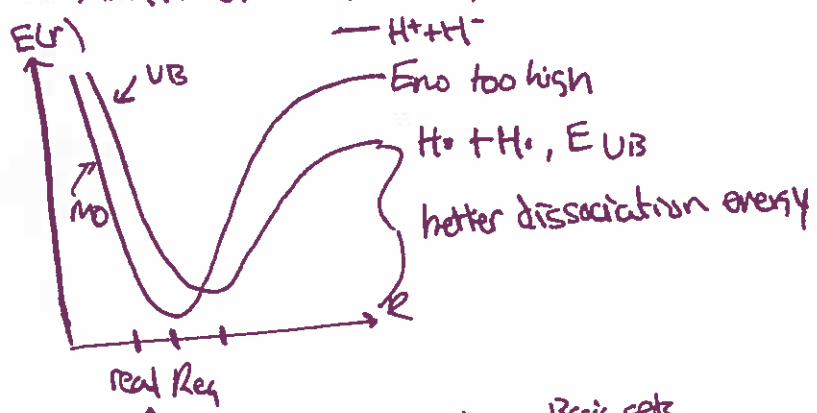
The delocalized MO does not dissociate correctly

Linus Pauling: keep only the 'good' terms

$$\psi_{VB} = C \{ |s_A\rangle_1 |s_B\rangle_2 + |s_B\rangle_1 |s_A\rangle_2 \} \{ \alpha_1\beta_2 - \alpha_2\beta_1 \}$$

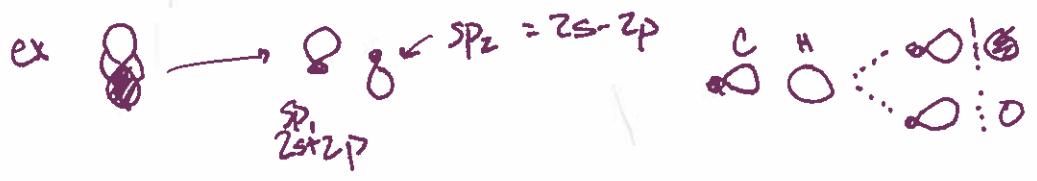
\uparrow
 valence bonds or hybrid basis function, always localized on each atom, not delocalized over the molecule

\Rightarrow better dissociation, worse equilibrium geometry



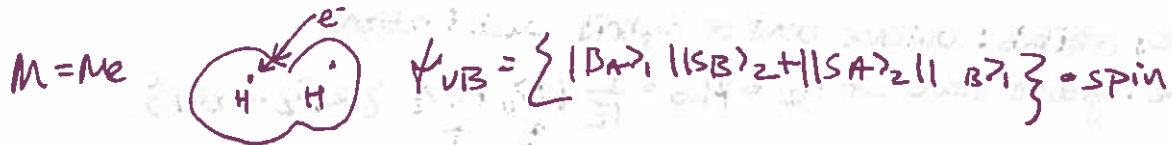
\uparrow experiment, or very large basis sets

We can construct all kinds of "hybrid" basis functions localized on a single atom



ex $\psi_{CH} \approx |\psi_1\rangle = \frac{1}{\sqrt{2}} \begin{vmatrix} |sp\rangle_1 \alpha_1 & |s_H\rangle_1 \beta_1 \dots \\ |sp\rangle_2 \alpha_2 & |s_H\rangle_2 \beta_2 \dots \end{vmatrix} = \psi_{VB}$

Interesting fact about valence bond basis functions



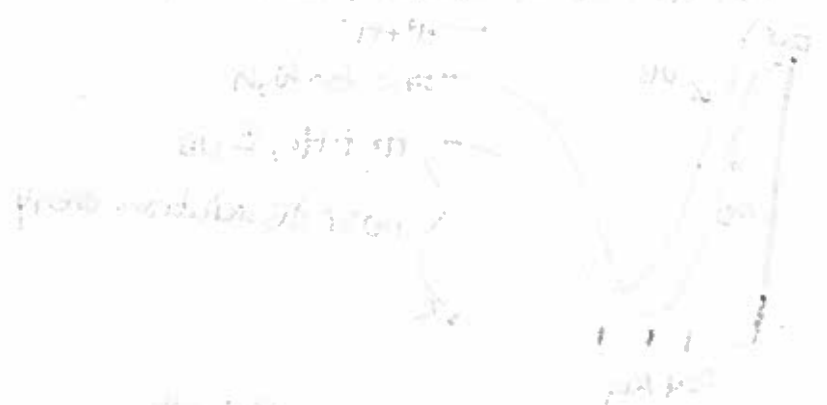
$\lim_{M \rightarrow \infty} \Psi_{UB} = \text{Lewis dot structure}$

$$\left\{ \frac{1}{\sqrt{2}}(|\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle), \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle), \frac{1}{\sqrt{2}}(|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \right\}$$

the more electrons the more the distribution of the bond of the spin up and down

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$$\frac{1}{\sqrt{2}}(|\uparrow\uparrow\rangle + |\downarrow\downarrow\rangle)$$