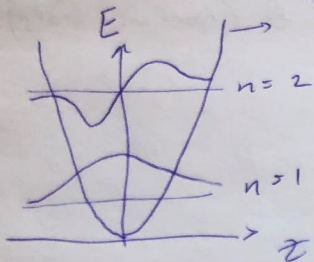


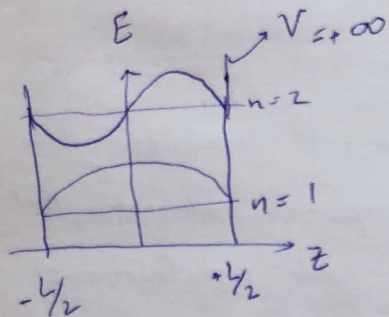
L10: review

$$V(z) = \frac{1}{2} k z^2$$



$$E_n = \hbar \sqrt{\frac{k}{m}} \left(n - \frac{1}{2} \right)$$

$n = 1, 2, \dots$



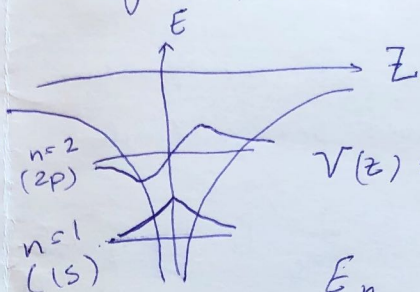
$$E_n = \frac{h^2 n^2}{8mL^2}$$

$n = 1, 2, \dots$

* the spring model ^{provides} a good approximation of diatomic bonds;

* the particle in box model provides a good approximation of energy states (and thus the absorption wavelength) of π electrons in conjugated systems.

today: H atom

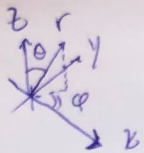


$$V(z) = \frac{e^2}{4\pi\epsilon_0} \cdot \frac{1}{z}$$

$$E_n \propto -\frac{1}{n^2}$$

* note that the dependence of $V(z)$ of H atom on z at high z values is much softer than that of spring model or box model, that is $V(z)$ does not change significantly by changing z at high $|z|$ values \Rightarrow energy of H atom changes much less steeply for higher quantum numbers ($E_n \propto -\frac{1}{n^2}$ for H atom vs $E_n \propto n$ for spring model)

The main difference today: 3D vs 1D



$$r = \sqrt{x^2 + y^2 + z^2}$$

* there is one quantum number per coordinate.

$$\text{CM: } E = \frac{1}{2m_e} (p_x^2 + p_y^2 + p_z^2) - \frac{e^2}{4\pi\epsilon_0 r}$$

$$\text{QM: } \hat{H} \psi = E_{nlm} \psi$$

$$-\frac{\hbar^2}{2m_e} \left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} \right) - \frac{e^2}{4\pi\epsilon_0 r}$$

the three quantum numbers are $n, l,$

$$\text{and } m_l: \begin{cases} n = 1, 2, 3, \dots \\ l = 0, 1, \dots, n-1 \\ m_l = -l, -l+1, \dots, l-1, l \end{cases}$$

Some examples of ψ_{nlm} :

$$\textcircled{1} \psi_{2,1,1} \sim re^{-r/a_0} \sin\theta (\cos\phi + i\sin\phi)$$

$\begin{matrix} \downarrow & \downarrow & \downarrow \\ n & l & m_l \end{matrix}$

$$\textcircled{2} \psi_{2,1,-1} \sim re^{-r/a_0} \sin\theta (\cos\phi - i\sin\phi)$$

$$\textcircled{3} \psi_{2,1,0} \sim re^{-r/a_0} \cos\theta$$

$$a_0 = \frac{4\pi\epsilon_0 \hbar^2}{m_e \cdot e^2} \approx 0.529 \text{ \AA} \text{ (Bohr length)}$$

! Plug in the above ψ examples in the Schrödinger equation and show that they are indeed valid eigenfunctions.

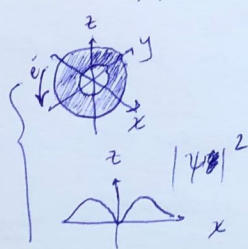
Now let's talk about how some of these functions look like...



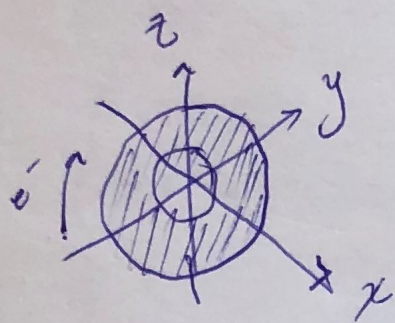
$$\textcircled{1} \psi_{2,1,1}$$

$$\cos\phi + i\sin\phi = e^{i\phi}$$

$$|\psi|^2 \sim |e^{i\phi}|^2 = 1$$



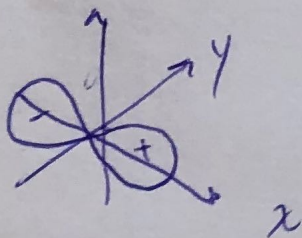
$$(2) \psi_{2,1,-1} \rightarrow m_l$$



m_l dictates the angular momentum of electron, i.e. the direction of electron's orbit.

the real (not imaginary) more convenient solutions of H model are formed by linear combination of the wavefunctions:

$$(1) + (2) : \psi_{p_x} \sim r e^{-r/a_0} \sin\theta \cos\varphi$$



$$(1) - (2) : \psi_{p_y} \sim r e^{-r/a_0} \sin\theta \sin\varphi$$

