

### Exam 1 ANSWER KEY

1. (10 pts) **Calculate** the energy separation (in Joules) between the  $n = 2$  and  $n = 3$  levels of a fluorine molecule confined in a one-dimensional box of length 1 cm. The energy levels for the particle in a box are given by

$$E_n = \frac{h^2 n^2}{8mL^2}$$

Answer: One fluorine atom is 18.998 amu, one fluorine molecule would then be 37.996 amu. One mole,  $6.626 \times 10^{23}$  molecules, of fluorine molecules would be 0.037 kg implying a single fluorine molecule is  $6.31 \times 10^{-26}$  kg (5 pts)

$$\Delta E = \frac{(6.626 \times 10^{-34})^2}{8(6.31 \times 10^{-26})(0.01)^2} (9 - 4) = 4.35 \times 10^{-38} \text{ J (5 pts)}$$

2. (10 pts) The Lennard-Jones potential for argon is given by the formula:

$$V(r) = 502.8 \left[ \left( \frac{3.34}{r} \right)^{12} - \left( \frac{3.34}{r} \right)^6 \right]$$

where  $r$  is in Ångströms. **Calculate** the force between two atoms,  $F(r) = -\partial V / \partial r$ . **Show** that  $F$  approaches 0 as  $r$  approaches  $\infty$ . **Does**  $F$  equal 0 at any other value of  $r$ ? If so, **at what value** of  $r$ ?

Answer:

$$F(r) = -\frac{\partial V}{\partial r} = -502.8 \left[ (-12) \left( \frac{3.34^{12}}{r^{13}} \right) - (-6) \left( \frac{3.34^6}{r^7} \right) \right] = 3016.8 \left[ 2 \left( \frac{3.34^{12}}{r^{13}} \right) - \left( \frac{3.34^6}{r^7} \right) \right] \text{ (5 pts)}$$

$$\lim_{r \rightarrow \infty} 2 \left( \frac{3.34^{12}}{r^{13}} \right) = 0 \text{ and } \lim_{r \rightarrow \infty} \left( \frac{3.34^6}{r^7} \right) = 0 \text{ implies that } \lim_{r \rightarrow \infty} F(r) = 0 \text{ (3 pts)}$$

$$F = 0 \text{ when } 2 \left( \frac{3.34^{12}}{r^{13}} \right) = \left( \frac{3.34^6}{r^7} \right) \text{ or } 2 \cdot 3.34^6 = r^6. \text{ Then } r = 3.34 \cdot 2^{1/6} \text{ \AA}.$$

at any other  $r$  value. As two argon atoms move further away, the force each exerts on the other decreases. (2 pts)

3. (15 pts) An electron has mass  $m_e \approx 9.1 \times 10^{-31}$  kg.

a. Its velocity has a range of  $\Delta v = 5$  m/s. What is the range of positions that will be measured?

b. In a helium atom, the range of velocities for an electron is closer to two million meters/second ( $2 \times 10^6$  m/s). What is the range of positions,  $\Delta x$ , that will be measured in meters? In Ångstroms? In nanometers?

c. How does the length in (b) compare to the Bohr radius:  $5.29 \times 10^{-11}$  m? What is the significance of that comparison?

Answer:

a. Starting with the Heisenberg uncertainty principle

$$\Delta x \Delta p = \frac{\hbar}{2}, \text{ we have the relationship } \Delta x = \frac{\hbar}{2\Delta p} \text{ and } \Delta p = m_e \Delta v.$$

If  $\Delta v = 5 \text{ m/s}$  and  $m_e \approx 9.1 \times 10^{-31} \text{ kg}$ ,  $\Delta p = 4.55 \times 10^{-30} \text{ kg}\cdot\text{m/s}$  and  $\Delta x = 1.16 \times 10^5 \text{ m}$ . (5 pts)

b. Same work as above, but this time  $\Delta v = 2.0 \times 10^6 \text{ m/s}$  making  $\Delta p = 1.82 \times 10^{-24} \text{ kg}\cdot\text{m/s}$  and  $\Delta x = 2.90 \times 10^{-11} \text{ meters} = 2.90 \times 10^{-1} \text{ \AA} = 2.90 \times 10^{-2} \text{ nm}$  (5 pts)

c. The length in (b) is a little more than half (0.547 factor difference) of the Bohr radius. This suggests that the most probable size of the 1s orbital of a helium atom is smaller than that of hydrogen. (5 pts)

4. (20 pts) Consider the sound wavefunction

$$\Psi(t) = \exp[-(t/2\Delta t)^2] \exp[i\omega t]$$

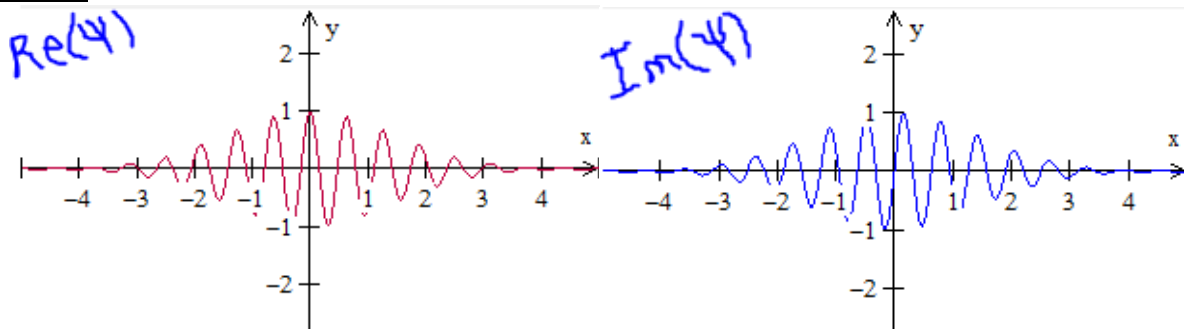
a. **Make two separate sketches** of the real and of the imaginary parts of  $\Psi(t)$  where  $\Delta t = 1 \text{ s}$  and  $\omega = 10 \text{ s}^{-1}$ , from  $t = -5$  to  $5 \text{ s}$ . Remember even/oddness of the sin and cos functions.

b. Consider only the real part of  $\Psi(t)$ , and **create another sketch** of what would happen if  $\Delta t$  was increased to  $2 \text{ s}$ .

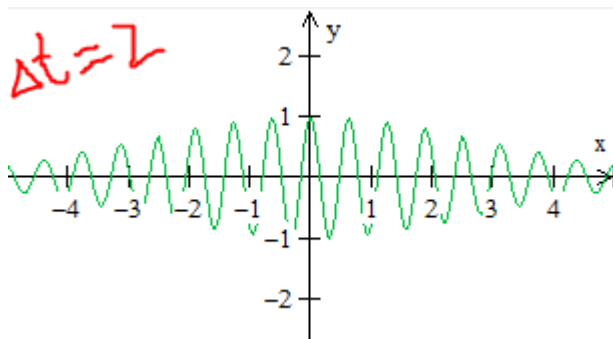
c. Now **sketch** what would happen to the real part if  $\Delta t$  remains  $1 \text{ s}$ , but the frequency is increased to  $\omega = 20 \text{ s}^{-1}$ .

d. **Does** taking the complex conjugate of  $\Psi(t)$  change either of your sketches from part (a)? If so, **how?** (You can state the answer in words – no sketch required.)

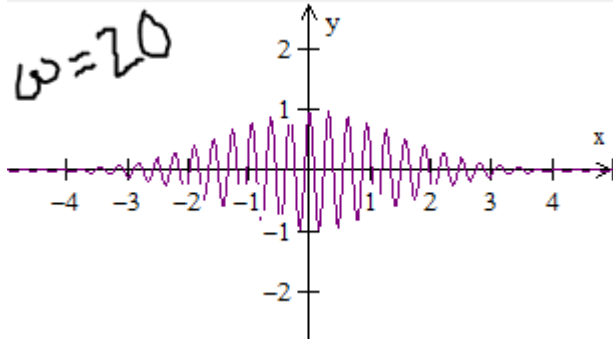
Answer:



a.  
(5 pts)



b. (5 pts)



c. (5 pts)

d. Taking the complex conjugate of the wavefunction only changes the imaginary part. The imaginary part of the wavefunction is now negated ( $\sin(x) \rightarrow -\sin(x)$ ). (5 pts)

5. (20 pts) Let a quantum particle be in state  $\Psi(x, t)$ . In this question, we will use postulate 4 to show step-by-step that the average value  $\langle A \rangle = \bar{A}$  of observable  $A$  is given by

$$\bar{A} = \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{A} \Psi(x, t),$$

where  $\hat{A}$  is the operator for observable  $A$ .

a. Recall from the \$ bills homework that  $\bar{A} = \sum a_n P(A = a_n)$ , and from postulate 4 that  $P(A = a_n) = \left| \int_{-\infty}^{\infty} dx \varphi_n^*(x) \Psi(x, t) \right|^2$ . **Combine these two equations** to write an expression for  $\bar{A}$  in terms of the eigenvalues and eigenfunctions of  $\hat{A}$ .

b. Now **write down** the same expression for  $\bar{A}$ , but multiply out the  $|\cdot|^2$  square modulus containing the integral explicitly. Something like

$$\int_{-\infty}^{\infty} dx \Psi^*(x, t) \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t)$$

should appear in your expression (the integral times its complex conjugate). It's always a good idea to give your integration variables different names (here  $x$  and  $x'$ ) when you have several integrals in an expression!

c. Your expression for  $\bar{A}$  has an “ $a_n$ ” outside the first integral. You can stick it inside the integral in front of  $\varphi_n(x)$ , and play the reverse of a trick we have done in class several times: since  $\hat{A} \varphi_n(x) = a_n \varphi_n(x)$ , you can replace “ $a_n$ ” **by what** in the integral? **Write down** the formula for  $\bar{A}$  again.

d. Remember that any function can be expanded as a linear combination of a complete set of functions. For example,  $\Psi(x, t) = \sum c_m \varphi_m(x)$ , where  $c_m = \int_{-\infty}^{\infty} dx' \varphi_m^*(x') \Psi(x', t)$  is the “overlap integral” between  $\varphi_m$  and  $\Psi$ . Look at your expression from c., and you should be able to **replace one of the integrals** by  $c_m$ , and then **do the sum**  $\sum c_m \varphi_m(x)$ , which equals  $\Psi(x, t)$ . If all went right, you should now have the expression

$$\bar{A} = \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{A} \Psi(x, t)$$

that you were supposed to prove. Congrats, in 1926 you would have won a Nobel Prize!

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Answer:

a. (5 pts)

$$\bar{A} = \sum_n a_n \left| \int_{-\infty}^{\infty} dx \varphi_n^*(x) \Psi(x, t) \right|^2$$

b. (5 pts)

$$\bar{A} = \sum_n a_n \int_{-\infty}^{\infty} dx \Psi^*(x, t) \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t)$$

c. (5 pts)

$$\begin{aligned} \bar{A} &= \sum_n \int_{-\infty}^{\infty} dx \Psi^*(x, t) a_n \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t) \\ &= \sum_n \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{A} \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t) \end{aligned}$$

d. (5 pts)

$$\begin{aligned} \bar{A} &= \sum_n \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{A} \varphi_n(x) \int_{-\infty}^{\infty} dx' \varphi_n^*(x') \Psi(x', t) \\ &= \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{A} \sum_n \varphi_n(x) \cdot c_n = \int_{-\infty}^{\infty} dx \Psi^*(x, t) \hat{A} \Psi(x, t) \end{aligned}$$

